Random laser from engineered nanostructures obtained by surface tension driven lithography

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Abstract. The random laser emission from the functionalized thienyl-S,S-dioxide quinquethiophene (T5OCx) in confined patterns with different shapes is demonstrated. Functional patterning of the light emitter organic material in well defined features is obtained by spontaneous molecular self-assembly guided by surface tension driven (STD) lithography. Such controlled supramolecular nano-aggregates act as scattering centers allowing the fabrication of one-component organic lasers with no external resonator and with desired shape and efficiency. Atomic force microscopy shows that different geometric pattern with different supramolecular organization obtained by the lithographic process tailors the coherent emission properties by controlling the distribution and the size of the random scatterers.

1.Introduction

Polymeric and low-molecular weight organic compounds are considered innovative and promising materials because they are efficient light emitters in a wide spectral range, biodegradable, chemically stable, low cost and easily usable for deposition on rigid and flexible substrates. Among them thiophene-based oligomers have shown relevant photoluminescence that makes them suitable candidates for the realization of photonic devices, such as organic light-emitting diodes (OLEDs) and solid state lasers.[1–3] Recently both amplified spontaneous emission and lasing properties of single crystals of thiophene oligomers have been reported, showing that the crystalline material acts as gain medium and the resonators are a pair of facets of the crystal,[4, 5] or microcavities of different shape fabricated by electron-beam lithography and dye etching,[6, 7] or distributed feedback cavities.[8]

These approaches allow for remarkable device performances but require complex technological steps and device architectures which are not easily integrable in practical applications. Recently promising methods for the fabrication of planar lasers, based on an active molecular layer in which defects, aggregates or external beads behave as scattering centers, have been reported.[9–12] This class of devices, namely random lasers, do not involve any additional external feedback structures, thus are simple to fabricate and low cost. However, due to the intrinsically randomness of the scattering centers, conventional methods for the fabrication of random lasers do not allow for a careful control of the device geometrical parameters, and in turn of the lasing properties.

Here we demonstrate the random laser emission from scattering nano-aggregates of a thiophene based molecule, obtained in a controlled way by a simple soft lithography technique.

Functional patterning of the organic material in well defined features is obtained by exploiting and controlling the instability phenomena affecting a liquid thin film on a non-wetting surface [2, 13, 14] and the intrinsic self-organization of π -conjugated molecules in spatially restricted environments.[1, 13] Such a bottom-up approach allows for both the modulation of the single molecule conformation and the supramolecular organization, thus permitting to finely control a range of physical parameters typical of confined systems, without any conventional, complex and expensive lithographic step.[15] Quasi-equilibrium conditions in which self-organization evolves are, in fact, strongly

affected by dimensionality and energetic parameters at the solid-liquid interface. Therefore, by modulating specific process parameters we succeed in finely tuning dimensions and spatial distribution of the functional nanostructures that, in turn, strongly affect the emission properties, as the lasing threshold.

As prototype active material we use a functionalized thienyl-S,S-dioxide quinquethiophene (T5OCx, structure in Fig. 1a), which has been deeply investigated as active material for lasing applications. Amplified spontaneous emission (ASE) from spin coated thin films, [16] as well as coherent lasing in external mirrors [17] or distributed feedback (DFB) based architectures [18] have been demonstrated for T5OCx. In our case no external cavity or conventional DFB structures are exploited, being the feedback provided by the random distribution of highly scattering centers between which light is entrapped. The difference with standard random lasers [19–21] is that in the latter the scattering centers are usually high refractive index colloidal particles or powders dispersed in the active dye matrix, as well as defects in the active films not easy to control in dimension, shape and distribution. Differently, our systems consist of scattering centers obtained by the accurate control of active molecules self-assembly. Such a controlled architecture allows for single peak lasing emission increasing in intensity and getting narrow in lineshape as pump energy increases, differently from what is generally observed in so far reported thiophene based materials, where sharp weak peaks superimposing the ASE spectrum appear. [16, 22] The single peak is typical of non resonant feedback random lasers [21] originated from extended modes over a substantial portion of the gain volume. To our knowledge, this is the first demonstration of collective lasing effect tuned by randomly distributed supramolecular assemblies realized by soft lithography.

2. Fabrication of surface tension driven patterns

The synthesis and characterization of 3,3',4"',3""tetracyclohexyl-3",4"-di(n-hexyl)- 2,2':5',2"':5"',2"':5"",2""-quinquethiophene-1",1"-dioxide (T5OCX) have already been reported.[23] For the preparation of T5OCX used in this study, an improved procedure was employed taking advantage of ultrasounds for the bromination reactions and microwaves for coupling reactions, using the typical conditions described in reference.[24]

Molecular micro-structures are realized by depositing few micro liters of 0.4% (w/v) solution of T5OCx in dichloromethane (CH_2Cl_2) on different stamps used as geometrical constrain. The molecular patterns are obtained by exploiting surface tension driven (STD) deposition under the effect both of energetic and geometric constrain, in which the amplification of the surface instabilities at the liquid interface, driven by an external template, realizes a controlled pattern of T5OCx. In Fig. 1a we show the 3D molecular structure of T5OCx and in Figs. 1b-d schematic illustration of three different bottom-up lithographic techniques:

Ring: A copper (Cu) disk (d=2 mm) is fixed on a glass substrate and used as template. The solution droplet is deposited on the top of Cu disk (Fig. 1b,e).[2, 13]

Stripe: The molecular stripes are carried out by MIMIC, exploiting microfluidic dynamics inside a poly(dimethylsiloxane) (PDMS) stamp to pattern soluble molecular materials. The stamp is realized by replica molding of a SU8 master. The microchannel (80 μm -wide stripe with a depth of 150 μm) is defined by placing the PDMS mold in conformal contact with a glass substrate (Fig. 1c,f).[25, 26]

Pixel microarrays: For the molecular pixels we used a metallic template fixed on a glass substrate. As a template, transmission electron microscopy (TEM) calibration Cu grid (G400, Gilder) with about 37 μm wide square holes and 25 μm wide bars was used (Fig. 1d,g).[2, 13]

Cartoons of the correspondent patterns are shown in Figs. 1e-g and the optical images in true colors in Figs. 1h-j show the final structures.

A fine modulation of the shape and distribution of the scattering centers is obtained for each pattern also by acting on the process kinetics. It is known, in fact, that differences in supramolecular packing due to self-organization of molecular material occurring during nucleation phenomena can be selectively addressed by controlling different type of energetic parameters (i.e., surface energy at interface, geometric confinement, process temperature etc).[2, 14]

The optical properties of the STD micro-patterns are investigated by confocal microscope (FluoView1000, Olympus) using a 40X objective (N.A. 0.85), a multi-line Ar laser (515 nm) and the corresponding dichroic mirror (DM 488 nm/515 nm). Figure 2 shows the confocal characterizations of the T5OCx stripe-shaped (see Figs. 2a,c,e,g) and micropixels (see Figs. 2b,d,f,h). Confocal x-z optical section (see Fig. 2c and d), 3D reconstruction of the organic active layers (see Fig. 2e and f) and the spatially resolved photoluminescence spectra (SR-PL) (see Fig. 2g and h) show that the STD lithographic process does not affect the optical properties of T5OCx.

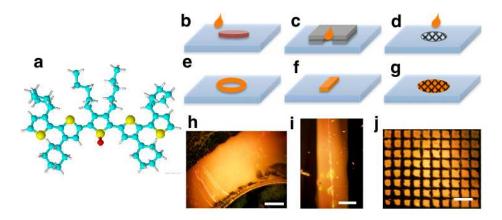


FIG. 1: (online color at: www.lpr-journal.org) Schematic illustration of the molecular material used and of the STD patterning procedures. a) 3D molecular structure of T5OCX. b-d) Sketches of the different bottom-up lithographic techniques adopted to pattern the thiophene dye: ring (b), stripe (c) and squared pixels (d). e-g) The corresponding cartoons of the realized patterns and h-j) optical images in true color of the fabricated samples. A drop of T5OCx solution in CH_2Cl_2 is poured at the centre of the related template placed on a glass substrate. The average thickness of all samples is about 1μ m. Scale bars: $100 \ \mu m$ (h,j) and $50 \ \mu m$ (i).

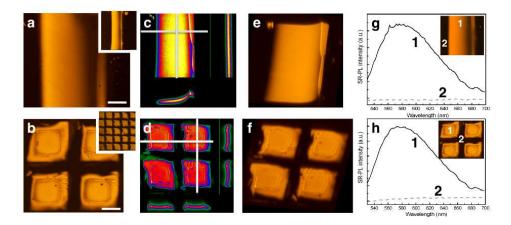


FIG. 2: (online color at: www.lpr-journal.org) (a,b) Laser scanning confocal microscopy images of the T5OCx stripe-shaped (a) and micropixels array (b) patterned by STD technique. Scale bars: $20~\mu m$; (c,d) z-projection in false color of confocal scanning in depth of stripe (c) and pixels (d) patterns with a resolution step of 200~nm for each slice. The x-z and y-z optical section are reported in the image sides; (e,f) 3D reconstruction of the spatial resolved photoluminescence emission, obtained by x-z optical sections of confocal scans; (g,h) spatially resolved photoluminescence spectra (SR-PL) of stripe-shaped (g) and micropixels array (h) acquired inside (1) and outside (2) the molecular pattern.

3.Random lasing from supramolecular nano-structures

The patterned samples are pumped by a frequency doubled Q-switched frequency doubled Nd:YAG pulsed laser emitting at λ =532 nm, with 10 Hz repetition rate, 8 ns pulse duration and with 8 mm beam diameter. For each sample, the laser spot is tuned by using an iris at the exit and the emitted radiation is collected from the input face of the sample at an angle of about 40° from the beam direction after focalization into an optic fiber connected to a spectrograph equipped with electrically cooled CCD array detector. A sketch of the experimental set-up is reported in Fig. 3a. We probe different emission directions in the angular range 25°-65° from the incident beam direction and we do not observe any significative variation in the emission profiles. The exposure time was set to 1 s and each spectrum results from the average of 10 pump shots.

In Fig. 3 the emission spectra from stripe (b) and ring (c) confinement geometries are reported. The amplified stimulated emission from both systems is clearly evidenced by the appearance over the broad gain band of a peak at 610 nm with a linewidth narrowing down to an order of magnitude as the pump energy is increased. In these

experiments the spot size of the incident beam is about 8 mm for the stripe and 2 mm for the ring, meaning that only a portion of the ring is illuminated. In order to improve the lasing properties of the fabricated random lasers we prepare two different samples of the ring laser, with initial better performances with respect to the stripe one, at different process temperatures, $T_{slow} = 24$ °C and $T_{fast} = 50$ °C, corresponding to slow and fast molecular assembly kinetics, respectively.

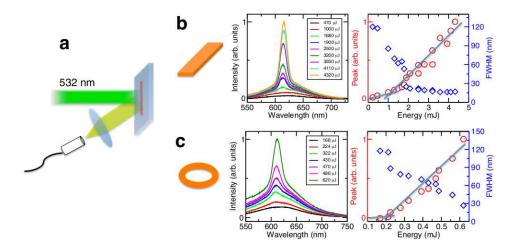


FIG. 3: (online color at: www.lpr-journal.org) (a) A sketch of the random laser set-up. (b,c) Emission spectra and corresponding spectral analyses of a portion of the stripe-shaped (b) and ring-shaped (c) dye self assemblies at variance with the injected pump energy. In the right panels the peak intensities (red) and the full width at half maximum (blue) are reported.

In Fig. 4 atomic force microscopy (AFM, Multimode 8-Bruker) characterization of the ring laser is reported, showing the conformational structure of the patterned organic dye at the nano scale. In particular, an amorphous and quite homogeneous arrangement (dark part) of T5OCx with the presence of randomly distributed disk-shaped nanostructures (bright spots) is evident in both samples. Size and distribution density of the aggregates instead strongly depend on the process temperature, with smaller aggregates and larger relative distances in the case of fast kinetics (see Fig. 4c and f). The formation of these well-defined self-assembled nanostructures, dispersed across the amorphous bulk, is induced in both samples by the quasi-equilibrium self-assembling dynamics under energetic confinements of the STD lithography. In particular, it is known that the modulation of molecular arrangement in the solid-state of a thiophene-based molecule, driven by a STD process, is due to a balance of polar and apolar contributes to surface energy. [14] Conversely, the conventional spin-coating deposition of T5OCx is exclusively characterized by an amorphous arrangement without any supramolecular structure. [16] The different dimension and distribution of the aggregates can be explained by considering that a slow kinetics, nearest to the quasi-equilibrium regime, makes easier the molecules to arrange in a crystalline packing.[1] Conversely, a fast dynamics enhances the propagation of instabilities at the liquid interface during the STD process, thus favoring the formation of more distorted amorphous phase. Interestingly in Fig. 4b and e the correspondent phase images exhibit additional spots, not visible from the height signal, showing that such observed aggregates are present even in the space below the observable amorphous bulk and thus they are three dimensional spatially distributed.

The emission from the two ring-shaped lasers taken in the same experimental conditions is investigated and the relating lasing peak intensities reported in Fig. 4g, where clear lasing thresholds of 70 μ J and 130 μ J are evident for T_{slow} and T_{fast} devices, respectively. Better lasing properties are obtained in T_{slow} device where more packed and bigger supramolecular nanostructures (bright spots) provide better feedback conditions. Statistical analysis of the AFM images reported in Table I show in fact noticeable difference in the aggregates density, size and relative distance.

The presented AFM investigation evidences that our molecular lasers are random lasers with the T5OCx aggregates acting as scattering centers. Emission spectra at higher resolution (0.3 nm) do not show Fabry-Perot like oscillations as those reported for single crystal thiophenes [4], confirming that our systems are random lasers. The single non-resonant random laser peak is thus explained by the fact that the pumped region of millimeter size is much larger than the average distance between the scatters of the order of micrometer and the single localized modes are highly interacting leading to one intense coherent peak. Such single peak is observed even in single shot measurements with 0.1 s pulse duration, giving non-averaged emission spectra. This occurs even though the samples are weakly scattering as indicated by the light transport mean path l_{tr} reported in Table I. These values are obtained, in the approximation

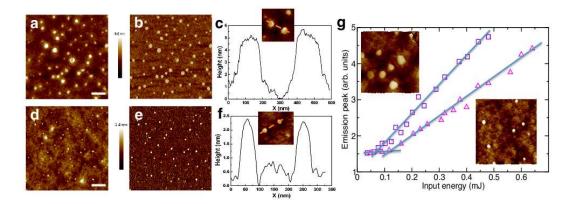


FIG. 4: (online color at: www.lpr-journal.org) a-c) AFM topography (a) and phase (b) images of the ring-shaped thiophene sample prepared at room temperature ($T_{slow} = 24^{\circ}C$) and section profile along the dashed lines of representative portions (in the insets) (c); d-f) AFM topography (d), phase (e) images and section profile (f) of the ring-shaped thiophene sample prepared at temperature of $T_{fast} = 50^{\circ}C$. Scale bars: 500 nm. The brigth spots indicate the T5OCx self-assembled nanostructures that act as scattering centers. g) Peak emission intensities of samples prepared at $T_{slow} = 24^{\circ}C$ (squares) and $T_{fast} = 50^{\circ}C$ (triangles); insets: correspondent AFM zooms.

TABLE I: Statistical analysis of the AFM images obtained by NanoScope software analysis over a $3\mu m \times 3\mu m$ frame and lasing energy treshold.

 N_{ns} is the surface number density of the scattering nanostructures; d is the average diameter of the nanostructures; H is the average height; A is surface area of all the nanostructures; L is the average distance between the nanostructures; R_q is the root mean squared roughness; l_{tr} is the estimated transport mean free path; I_t is the lasing threshold intensity, calculated as the energy threshold divided by the illuminated surface area and the pump pulse duration.

Sample	$N_{ns} (\mu m^{-2})$	d (nm)	H (nm)	$A (nm^2)$	L (nm)	R_q (nm)	$l_{tr}(\mu m)$	$I_t(\mathrm{J/m}^2\mathrm{s})$
Ring (T_{slow})	8.7	83 ± 50	3.5 ± 1.2	7.35×10^{3}	500 ± 300	0.950	16	3.7×10^{9}
Ring (T_{fast})	6.6	$55\pm\ 34$	$1.7 \pm\ 0.5$	$3.17{\times}10^3$	$600 {\pm} 450$	0.353	108	6.9×10^{9}
Micropixel	5.1	$48 \!\pm 18$	1.9 ± 0.8	1.95×10^{3}	1260 ± 800	0.718	250	2.5×10^{14}

of weakly scattering regime, by considering $l_{tr} = 1/(\sigma_{tr}N_{ns})$, with σ_{tr} the transport cross section and N_{ns} the number density of the scattering nanostructures as listed in Table I. Being the propagation of scattered light in our samples mainly in the 2 dimensions perpendicular to the beam propagation, we approximate our scattering nanosctructures to cylinders infinitely long along the beam direction and with diameter d, thus we use the analytical expression of the transport cross section σ_{tr} as in [16, 27], with the index of refraction n=1.8 as measured and reported previously [16].

Moreover, we consider several samples and find that only those with certain amount of disorder sustain random lasing action. As an example the pixel array device in Fig. 1j does not show any lasing with ns-pulse input energy, due to the fact that the internal region presents supramolecular dye assemblies similar to the other samples, but with visible lower distribution and reduced dimensions, as reported in Table I and shown in the AFM images of Fig. 5b-d. This negatively affects the feedback mechanism and amplified emission is observed only at much higher peak power obtained by using for optical pumping a multi-stage frequency doubled Ti:Sapphire laser system (Coherent Hidra) with 10 Hz repetition-rate, 50 fs pulse duration and wavelength λ =400 nm. In Fig. 5e the onset of the peak by increasing pumping energy evidences the random laser emission. In these experiments the spot size of the incident beam is about 5 mm meaning that the emission is from a group of pixels as in the inset of Fig. 2b. The low efficiency compared to the results in Fig. 3 is due i) to the fact that the free standing pixels do not fill the entire illuminated surface; ii) to the reduced size and distribution of the scattering aggregates for each pixel as indicated by the section profile along the dashed line depicted in Fig. 5d and by the values in Table I.

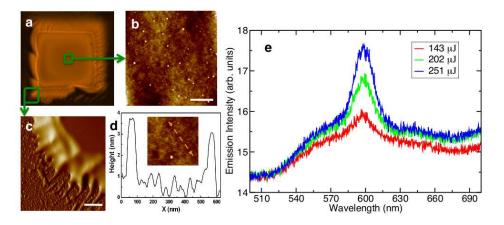


FIG. 5: (online color at: www.lpr-journal.org) a-c) Confocal image of one single pixel (a) and AFM images of two portions (b,c). Scale bars: 500 nm (b); 1 μ m (c). d) Section profile along the dashed line of a portion of AFM image from the internal part of a single pixel shown in the inset. e) Emission spectra of the ensemble of pixels pumped by an amplified 50fs λ =400nm Ti:Saphire laser with spot size 5 mm.

4. Conclusion

In conclusion, the lasing properties of supramolecular self assemblies of a suitable thiophene-based organic dye is investigated. The use of surface tension driven (STD) lithographic processes allows to obtain organic mini-lasers of different shapes and importantly to tailor the structure of the random lasers at the nanoscale by finely tuning the supramolecular self assembly of the organic dye. The peculiarity of the STD structures, capable of being greatly modulated, combined with the intrinsic conformational flexibility of several substituted oligothiophenes, [28] give the possibility to a priori select shape and performance characteristics of the final device. By using atomic force microscopy it is shown that the coherent radiation emission is a random laser where the scattering centers are thiophene aggregates formed by spontaneous molecular self-assembly guided by STD lithography. The optimization of the deposition procedure and process kinetics lead to tailor the coherent emission properties by controlling the distribution and the size of the random scatterers. Our results open the way to the fabrication of one-component organic lasers with no external resonator and with desired shape and efficiency by using simple soft lithographic techniques.

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